

6. "Anomalies"

[Peskin/Schroeder §19; Zee §IV.7]

→ Q.: symmetry of classical physics $\hat{=}$ symmetry of quantum physics?
 (action $S[\phi]$ invariant under $\phi \rightarrow \phi + \delta\phi$) (path integral $\int \mathcal{D}\phi e^{iS[\phi]}$ invariant)

→ A.: not always; measure $\mathcal{D}\phi$ may or may not be invariant.

example: rotational invariance of $Q\pi$
 would be strange if quantum fluctuations preferred a specific direction!

but: quantum fluctuation can break (some) classical symmetries;
 this phenomenon is called "anomaly";
 conceptually clear: change of integration variables \rightarrow don't forget Jacobian

important subject in QFT \rightarrow many ways of looking at it
 here: see a class. symmetry vanishing, by explicit Feynman diag calc.

- theory of one massless fermion: $\mathcal{L}_1 \equiv \bar{\psi} i \gamma^\mu \partial_\mu \psi$
 invariant under $\psi \rightarrow e^{i\theta} \psi$ and $\bar{\psi} \rightarrow e^{-i\theta} \bar{\psi}$
 conserved current $J^\mu \equiv \bar{\psi} \gamma^\mu \psi$ (vector) and $J_5^\mu \equiv \bar{\psi} \gamma^\mu \gamma^5 \psi$ (axial)

((check: $\partial_\mu J^\mu = 0 = \partial_\mu J_5^\mu$ via class. eqn. of motion $i \gamma^\mu \partial_\mu \psi = 0$))

- calculate (Fourier transform of) amplitude $\langle 0 | T J_5^\lambda(0) J^\mu(x_1) J^\nu(x_2) | 0 \rangle$

$$\Delta^{\lambda\mu\nu}(k_1, k_2) = \int \frac{d^4 p}{(2\pi)^4} \text{tr} \left(\gamma^\lambda \gamma^5 \frac{i}{\not{p} - \not{k}_1} \gamma^\nu \frac{i}{\not{p} - \not{k}_2} \gamma^\mu \frac{i}{\not{p}} + \gamma^\lambda \gamma^5 \frac{i}{\not{p}} \gamma^\mu \frac{i}{\not{p} - \not{k}_2} \gamma^\nu \frac{i}{\not{p} - \not{k}_1} \right)$$

if (in QFT) $\partial_\mu J^\mu = 0$ holds, then $k_{1\mu} \Delta^{\lambda\mu\nu} = 0$ and $k_{2\nu} \Delta^{\lambda\mu\nu} = 0$

\therefore $\partial_\mu J_5^\mu = 0$ holds, then $q_\lambda \Delta^{\lambda\mu\nu} = 0$

\rightarrow can check this by explicit computation.

6.1 Vector current conservation

- would non-conservation of either current be a disaster?

- J^μ : charge $Q = \int d^3x J^0$ counts # of fermions
 → would be difficult to interpret if not conserved!
 couple photon to ψ , photon line coming into vertex γ^μ
 would have propagator $\frac{-i}{k^2} (\partial_{\mu\nu} - (1-\xi) \frac{k_\mu k_\nu}{k^2})$
 → gauge dependence falls out if $b_{\mu\nu} \Delta^{\mu\nu} = 0$

- J^5 : who cares if axial charge $Q^5 = \int d^3x J^0_5$ changes in time?

- naive calculation

$$b_{\mu\nu} \Delta^{\mu\nu}(b_1, b_2) = i \int \frac{d^4p}{(2\pi)^4} \text{tr} \left(\cancel{\gamma^\lambda \gamma^5} \frac{1}{\cancel{p-q}} \cancel{\gamma^\nu} \frac{1}{\cancel{p-k_1}} \cancel{\gamma^\lambda} \frac{1}{\cancel{p}} + \cancel{\gamma^\lambda \gamma^5} \frac{1}{\cancel{p-q}} \cancel{\gamma^\nu} \frac{1}{\cancel{p-k_2}} \cancel{\gamma^\lambda} \frac{1}{\cancel{p}} \right)$$

$$= i \int \frac{d^4p}{(2\pi)^4} \text{tr} \left(\cancel{\gamma^\lambda \gamma^5} \frac{1}{\cancel{p-q}} \cancel{\gamma^\nu} \frac{1}{\cancel{p-k_1}} - \cancel{\gamma^\lambda \gamma^5} \frac{1}{\cancel{p-q}} \cancel{\gamma^\nu} \frac{1}{\cancel{p-k_2}} \right)$$

shift $p \rightarrow p - b_1$

$$= 0$$

- more careful calculation

when is it OK to shift integration variables?

1dim: $\int_{-\infty}^{\infty} dp [f(p+a) - f(p)] = \int_{-\infty}^{\infty} dp [a \partial_p f(p) + \mathcal{O}(a^2)] = a(f(\infty) - f(-\infty)) + \mathcal{O}(a^2)$

d dim: $\int d^d p_\epsilon [f(p+a) - f(p)] = \int d^d p_\epsilon [a^\mu \partial_{p^\mu} f(p) + \dots] \stackrel{\text{Gauss}}{=} \lim_{R \rightarrow \infty} \left\langle a^\mu \left(\frac{p^\mu}{R} \right) f(R) \underbrace{S_{d-1}(R)}_{\substack{\text{"surface" of } d \text{ dim sphere, radius } R}} \right\rangle$

for our 4 dim Nambu-Gim integral, from Wick rot.

$$\int d^4 p [f(p+a) - f(p)] = \lim_{R \rightarrow \infty} \left\langle i a^\mu \left(\frac{p^\mu}{R} \right) f(R) (2\pi^2 R^3) \right\rangle$$

angular average

use this for $a = -b_1$

and $f(p) = \text{tr} \left(\cancel{\gamma^\lambda \gamma^5} \frac{1}{\cancel{p-k_2}} \cancel{\gamma^\nu} \frac{1}{\cancel{p}} \right) = \frac{\text{tr}(\gamma^5 (\cancel{p-k_2}) \cancel{\gamma^\nu} \cancel{p} \gamma^\lambda)}{(p-k_2)^2 p^2} \stackrel{\text{trace}}{=} \frac{4i \epsilon^{\tau\nu\sigma\lambda} b_{2\tau} p_\sigma}{(p-k_2)^2 p^2}$

$$\Rightarrow b_{\mu\nu} \Delta^{\mu\nu}(b_1, b_2) = \frac{i}{(2\pi)^4} \lim_{R \rightarrow \infty} \left\langle i (-b_1)^\mu \frac{p^\mu}{R} \frac{4i \epsilon^{\tau\nu\sigma\lambda} b_{2\tau} p_\sigma}{(R-k_2)^2 R^2} 2\pi^2 R^3 \right\rangle$$

$$\left(\langle \frac{p_\mu p_\nu}{R^2} \rangle = \frac{2\nu\sigma}{4} R^2 \right) \stackrel{\text{angular avg}}{\downarrow} = \frac{i}{8\pi^2} b_1^\mu \epsilon^{\tau\nu\sigma\lambda} g_{\mu\sigma} b_{2\tau} = \frac{\epsilon}{8\pi^2} \epsilon^{\lambda\nu\tau\sigma} b_{1\tau} b_{2\sigma}$$

→ $b_{\mu\nu} \Delta^{\mu\nu} \neq 0$?! fermion # not conserved, we disintegrate

$\text{tr}(\gamma^\mu \gamma^\nu \gamma^5) = -4i \epsilon^{\mu\nu\sigma\lambda}$