

• important integrals

$$\Gamma(n) = (n-1)! = \int_0^\infty dt \, t^{n-1} e^{-t}$$

$$\int_0^\infty dx \frac{x^{2a-1}}{(x^2+1)^b} = \frac{\Gamma(a) \Gamma(b-a)}{2 \Gamma(b)} \quad (\text{if } \operatorname{Re}[a] > 0, \operatorname{Re}[b-a] > 0)$$

$$\frac{1}{A_1^{a_1} A_2^{a_2} \dots A_n^{a_n}} = \frac{\Gamma(a_1 + \dots + a_n)}{\Gamma(a_1) \dots \Gamma(a_n)} \int_0^1 dx_1 \dots dx_n \frac{x_1^{a_1-1} \dots x_n^{a_n-1} \delta(1 - \sum_{i=1}^n x_i)}{[x_1 A_1 + \dots + x_n A_n]^{a_1 + \dots + a_n}}$$

Feynman
parametrization

3.1 one-loop divergences in QCD

→ goal: evaluate 1-loop gluon self-energy diagrams

$$\text{Feynman diagram: } \text{gluon self-energy} = \text{ghost loop} + \text{gluon loop} + \text{quark loop} + \text{quark loop with cross}$$

Use dimensional regularization, $d^4k \rightarrow d^d k$, $d = 4 - 2\epsilon$

use Feynman gauge $\xi=1$ for gluon propagator (for simplicity)

• 1st diagram: consider one quark flavor, do $\sum_f$ in the end

$$\text{Feynman diagram: quark loop} = - \operatorname{Tr} \int \frac{d^d k}{(2\pi)^d} \left(i g \gamma^\mu T^a \right) \frac{i(\not{k} + m)}{k^2 - m^2 + i\epsilon} \left(i g \gamma^\nu T^a \right) \frac{i(\not{k} + \not{q} + m)}{(k+q)^2 - m^2 + i\epsilon}$$

↑ trace over γ 's and T 's
one closed fermion loop

Remember: Dirac matrices

$$\text{def } \{ \gamma^\mu, \gamma^\nu \} = 2 g^{\mu\nu} \mathbb{1} \quad ; \quad (\gamma^\mu)^\dagger = \gamma^0 \gamma^\mu \gamma^0$$

$$\operatorname{tr}(\mathbb{1}) = N \quad (\text{we take } N=4; \text{ also } N=2^{d/2}, N=d \text{ seen})$$

$$\operatorname{tr}(\gamma^\mu) = 0 \quad ; \quad \operatorname{tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) = 0 \quad ; \quad \dots$$

$$\operatorname{tr}(\gamma^\mu \gamma^\nu) = N g^{\mu\nu} \quad ; \quad \operatorname{tr}(\gamma^{\mu\nu\rho\sigma}) = N (g^{\mu\nu} g^{\rho\sigma} - g^{\mu\rho} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\rho}) \quad ; \quad \dots$$

$$= - g^2 (\operatorname{tr}(\mathbb{1})) (\operatorname{tr}(T^a T^a)) \int \frac{d^d k}{(2\pi)^d} \frac{N^{1/2}(k, k+q)}{(k^2 - m^2 + i\epsilon)((k+q)^2 - m^2 + i\epsilon)}$$

$$\begin{aligned} N^{1/2}(k, p) &\stackrel{\operatorname{tr}(\gamma^i \gamma^j)}{=} m^2 g^{13} + m \cdot 0 + (g^{12} g^{34} - g^{13} g^{24} + g^{14} g^{23}) k_2 p_4 \\ &= (m^2 - k \cdot p) g^{13} + k'_1 p'_3 + p'_1 k'_3 \end{aligned}$$

for denominator, use Feynman parameters, see p. 30

$$\frac{1}{(k^2 - m^2 + i\epsilon)((k+q)^2 - m^2 + i\epsilon)} = \int_0^1 dx \frac{1}{[(1-x)(k^2 - m^2 + i\epsilon) + x(k^2 + 2kq + q^2 - m^2 + i\epsilon)]^2}$$

$$= \int_0^1 dx \frac{1}{[(k+xq)^2 + \underbrace{x(1-x)q^2 - m^2 + i\epsilon}_{=-4}]^2}$$

now, shift $k \rightarrow k - xq$ under integral $\int d^d k$

$$= -g^2 \cdot 4 \cdot \frac{1}{2} \delta^{a_1 a_2} \int \frac{d^d k}{(2\pi)^d} \int_0^1 dx \frac{N^{1/2}(k-xq, k+(1-x)q)}{[k^2 - 4 + i\epsilon]^2}$$

numerator = $(m^2 - (k^2 - x(1-x)q^2 + \cancel{O(k^2)})) g^{1/2} + 2\cancel{k^{1/2} k^{1/2}} - 2x(1-x) q^{1/2} q^{1/2} + \cancel{O(k^2)}$

linear (in k) terms integrate to 0; $\rightarrow \frac{g^{1/2} k^2}{d}$, see p. 29 bottom

$$= (m^2 + (\frac{2}{d}-1)k^2 + x(1-x)q^2) g^{1/2} - 2x(1-x) q^{1/2} q^{1/2}$$

to evaluate the $\int d^d k$ integral, perform a Wick rotation $k^0 = i k_E^0$

$$\sqrt{k^2 - m^2 + i\epsilon} = k_0^2 - \vec{k}^2 - m^2 + i\epsilon = (k_0 + \sqrt{\vec{k}^2 + m^2 - i\epsilon})(k_0 - \Gamma + i\epsilon)$$

$\xrightarrow{\Gamma+i\epsilon} x \rightarrow \frac{k_0}{x\Gamma+i\epsilon}$, $0 = \sqrt{\frac{k_0^2}{x^2} - \vec{k}^2 - m^2 + i\epsilon}$, $\rightarrow = -\downarrow = +\uparrow : k^0 = i k_E^0$

$$= -g^2 \cdot 4 \cdot \frac{1}{2} \delta^{a_1 a_2} \int_0^1 dx \cdot i \left\{ [(m^2 + x(1-x)q^2) g^{1/2} - 2x(1-x) q^{1/2} q^{1/2}] I_2^0(4) + (1-\frac{2}{d}) g^{1/2} I_2^1(4) \right\}$$

where $I_n^a(4) \equiv \int \frac{d^d k_E}{(2\pi)^d} \frac{(k_E^2)^a}{[k_E^2 + 4]^n}$ (basic 1-loop "tadpole" integral)

$$= \frac{1}{(2\pi)^d} \frac{2\pi^{d/2}}{\Gamma(d/2)} \int_0^\infty dk_E \frac{k_E^{d+2a-1}}{(k_E^2 + 4)^n} |4|^{d/2+a-n}$$

$$= \frac{\Gamma(a+\frac{d}{2}) \Gamma(n-a-\frac{d}{2})}{(4\pi)^{d/2} \Gamma(d/2) \Gamma(n)} |4|^{n-a-d/2} \quad \text{see p. 30}$$

note that $I_2^1(4) = I_2^0(4) \cdot |4| \cdot \frac{d}{2} \frac{1}{1-\frac{d}{2}} = I_2^0(4) \frac{|4|}{\frac{d}{2}-1}$ ($x\Gamma(x) = \Gamma(x+1)$)

$$= -g^2 \cdot 4 \cdot \frac{1}{2} \delta^{a_1 a_2} \int_0^1 dx \cdot i \cdot I_2^0(4) \left\{ (m^2 + x(1-x)q^2 - |4|) g^{1/2} - 2x(1-x) q^{1/2} q^{1/2} \right\}$$

$\Delta > 0 : = 2x(1-x)q^2$

$$= -4ig^2 \delta^{a_1 a_2} (g^{1/2} q^2 - q^{1/2} q^{1/2}) \int_0^1 dx I_2^0(m^2 - x(1-x)q^2) x(1-x)$$

$$= \frac{\Gamma(2-\frac{d}{2})}{(4\pi)^{d/2}} \int_0^1 dx \frac{x(1-x)}{[m^2 - x(1-x)q^2]^{2-\frac{d}{2}}}$$