

# Curvilinear coordinates


# Curvilinear coordinates

- Changing coordinates
- Covariant differentiation of vector and tensor fields
- Application (next week): Milne coordinates

Recipes, without proof nor attempt at rigor!

# Motivation: an issue with polar coordinates

Consider the (2D) plane with

● Cartesian coordinates  $(x, y) \stackrel{\text{def}}{=} (x^{1'}, x^{2'})$  with associated basis vectors  $(\vec{e}_x, \vec{e}_y) \stackrel{\text{def}}{=} (\vec{e}_{1'}, \vec{e}_{2'})$  – note the at first bizarre primed indices as well as their positions;  the reference!

● the usual polar coordinates  $(r, \theta) \stackrel{\text{def}}{=} (x^1, x^2)$  with the associated basis vectors  $(\vec{e}_r, \vec{e}_\theta) \stackrel{\text{def}}{=} (\vec{e}_1, \vec{e}_2)$  – here with unprimed indices.

The position  $\vec{R}$  of a point P of the plane may be denoted

$$\vec{R} = x \vec{e}_x + y \vec{e}_y = \sum_{i'} x^{i'} \vec{e}_{i'} = x^{i'} \vec{e}_{i'}$$

But you know that you would NOT write  $\vec{R} = r \vec{e}_r + \theta \vec{e}_\theta = x^i \vec{e}_i$ . 

# Motivation: an issue with polar coordinates

If you look at two neighboring points  $P$  and  $P + dP$  and call  $d\vec{R}$  their separation, you will also write (Cartesian coordinates)

$$d\vec{R} = dx \vec{e}_x + dy \vec{e}_y = \sum_{i'} dx^{i'} \vec{e}_{i'} = dx^{i'} \vec{e}_{i'}$$

And perhaps are you then ready to write (in polar coordinates)

$$d\vec{R} = dr \vec{e}_r + d\theta \vec{e}_\theta$$

so that at least one formula:

$$d\vec{R} = dx^i \vec{e}_i$$

(1)

is valid in every coordinate system.

Fine? Let us see where this is leading to...

# Separation vector between neighboring points in arbitrary coordinates

It is tempting to write  $d\vec{R} = \frac{\partial \vec{R}}{\partial x^i} dx^i$ .

An issue is that we know how to express  $\vec{R}$  as a function of the Cartesian coordinates  $\{x^{i'}\}$ , but in general not in terms of arbitrary coordinates  $\{x^i\}$ .

But... we generally know how to express the  $\{x^i\}$  in terms of the  $\{x^{i'}\}$  and reciprocally: functions  $x^i(\{x^{i'}\})$  and  $x^{i'}(\{x^i\})$ .

For instance:  $r(x, y) = \sqrt{x^2 + y^2}$ ,  $\theta(x, y) = \arctan(y/x)$

and  $x(r, \theta) = r \cos \theta$ ,  $y(r, \theta) = r \sin \theta$

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👉 We may write  $\vec{R}(\{x^i\}) = x(\{x^i\})\vec{e}_x + y(\{x^i\})\vec{e}_y = x^{i'}(\{x^i\})\vec{e}_{i'}$

and now differentiate w.r.t.  $x^i$ :

$$\frac{\partial \vec{R}}{\partial x^i} = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} \quad \Rightarrow \quad d\vec{R} = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} dx^i \quad (3)$$

# Basis vectors of an arbitrary coordinate system

Identifying Eqs. (1) and (3), we find the “proper” basis vectors  $\{\vec{e}_i\}$ :

$$\left. \begin{aligned} d\vec{R} &= dx^i \vec{e}_i \\ d\vec{R} &= \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} dx^i \end{aligned} \right\} \vec{e}_i = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} \quad (4)$$

For the polar coordinates:  $x(r, \theta) = r \cos \theta$  ,  $y(r, \theta) = r \sin \theta$

$$\frac{\partial x}{\partial r} = \cos \theta , \quad \frac{\partial y}{\partial r} = \sin \theta \quad \Rightarrow \quad \vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$$

$$\frac{\partial x}{\partial \theta} = -r \sin \theta , \quad \frac{\partial y}{\partial \theta} = r \cos \theta \quad \Rightarrow \quad \vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$$

Well  $\vec{e}_r$  is fine and I knew it, but... this  $\vec{e}_\theta$  is not normalized to 1! 🤔

# Motivation: a problem with polar coordinates

For the polar coordinates:  $x(r, \theta) = r \cos \theta$  ,  $y(r, \theta) = r \sin \theta$

$$\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y \qquad \vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$$

Well  $\vec{e}_r$  is fine and I knew it, but... this  $\vec{e}_\theta$  is not normalized to 1! 🤔

Indeed! But who said it should? 🧐

Now  $\vec{e}_\theta$  is well-behaved in the limit  $r \rightarrow 0$ , when  $\theta$  is not defined: isn't this a nice feature, worth discarding the normalization? 😇

(the perplexed students): OK, we'll live with it... 😐

# Metric tensor

Start from eq.(1):  $d\vec{R} = dx^i \vec{e}_i$  and square it:

$$d\vec{R}^2 = (dx^i \vec{e}_i) \cdot (dx^j \vec{e}_j) = (\vec{e}_i \cdot \vec{e}_j) dx^i dx^j$$

Define  $g_{ij} \equiv \vec{e}_i \cdot \vec{e}_j$  and replace:  $d\vec{R}^2 = g_{ij} dx^i dx^j$

● Cartesian coordinates:  $g_{i'j'} = \delta_{i'j'} \Leftrightarrow d\vec{R}^2 = (dx)^2 + (dy)^2$

● Polar coordinates:\*  $g_{11} = g_{rr} = 1$  ,  $g_{22} = g_{\theta\theta} = r^2$  ,  $g_{r\theta} = g_{\theta r} = 0$   
 $\Leftrightarrow d\vec{R}^2 = (dr)^2 + r^2(d\theta)^2$

looks familiar...

$g_{ij}$  is called **metric tensor**.

(More precisely, these are its components, but let's not be picky today).

\* Using the expressions of  $\vec{e}_r$ ,  $\vec{e}_\theta$  in terms of  $\vec{e}_x$ ,  $\vec{e}_y$ , since we know how to compute the scalar product in Cartesian coordinates.

# Metric tensor and its inverse

Define  $g^{ij}$  such that its “product” (more accurately: contraction) with the metric tensor equals the identity:

$$g^{ij} g_{jk} = \delta_k^i$$

$g^{ij}$  is called the inverse metric tensor.

Cartesian coordinates:  $g^{i'j'} = \delta^{i'j'}$ ; polar coordinates:  $g^{rr} = 1$ ,  $g^{\theta\theta} = \frac{1}{r^2}$ .

$g_{ij}$  and  $g^{ij}$  can be used to lower / raise indices:

- No impact on Cartesian components.

- In polar coordinates: no impact on  $r$ -components, but  $\theta$ -components vary.

👉 What does it mean? Not much today...

# Next step: generalize!

Starting from coordinates  $\{x^{i'}\}$  with orthogonal\* and normalized\* basis vectors  $\{\vec{e}_{i'}\}$ , one goes to alternative coordinates  $\{x^i\}$ :

- with the basis vectors  $\vec{e}_i = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'}$  such that  $d\vec{R} = dx^i \vec{e}_i$  holds
- and the metric tensor  $g_{ij} \equiv \vec{e}_i \cdot \vec{e}_j$  such that  $d\vec{R}^2 = g_{ij} dx^i dx^j$ .

We may do that:

- ... in 2-dimensional Euclidean space:  $g_{i'j'} = \text{diag}(+1, +1)$
- ... in 3-dimensional Euclidean space:  $g_{i'j'} = \text{diag}(+1, +1, +1)$
- ... in 4-dimensional Minkowski space:  $g_{i'j'} = \text{diag}(-1, +1, +1, +1)$

\* Using the appropriate (pseudo)scalar product encoded in  $g_{i'j'}$

# Second motivation: another issue with polar coordinates

How should one differentiate?

Consider a (smooth enough) function  $f$ .

Its derivatives are supposed to measure the variations of the function in the various respective directions:

$$f(x + \delta x, y) \simeq f(x, y) + \frac{\partial f}{\partial x} \delta x$$

and so on...

# Second motivation: another issue with polar coordinates

How should one differentiate?

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$$f(r, \theta + \delta\theta) \simeq f(r, \theta) + \frac{\partial f}{\partial \theta} \delta\theta$$

and so on...

# Second motivation: another issue with polar coordinates

For instance:  $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

with (slide 6)  $\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$  and  $\vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$

Let us differentiate:

$$\frac{\partial \vec{f}(r, \theta)}{\partial \theta} = \frac{\partial \cos \theta}{\partial \theta} \vec{e}_r - \frac{1}{r} \frac{\partial \sin \theta}{\partial \theta} \vec{e}_\theta = -\sin \theta \vec{e}_r - \frac{\cos \theta}{r} \vec{e}_\theta$$

OK, fine!

Wait!  $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

$$= \cos \theta (\cos \theta \vec{e}_x + \sin \theta \vec{e}_y) - \frac{\sin \theta}{r} (-r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y)$$
$$= \vec{e}_x = \overrightarrow{\text{constant!}} \text{ How can its derivative be } \neq \vec{0}?$$

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For instance:  $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

with (slide 6)  $\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$  and  $\vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$

Let us differentiate: are also functions of  $(r, \theta)$ !

~~$$\frac{\partial \vec{f}(r, \theta)}{\partial \theta} = \frac{\partial \cos \theta}{\partial \theta} \vec{e}_r - \frac{1}{r} \frac{\partial \sin \theta}{\partial \theta} \vec{e}_\theta = -\sin \theta \vec{e}_r - \frac{\cos \theta}{r} \vec{e}_\theta$$~~

Use the product rule!

OK, fine!

Wait!  $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

$$= \cos \theta (\cos \theta \vec{e}_x + \sin \theta \vec{e}_y) - \frac{\sin \theta}{r} (-r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y)$$

$$= \vec{e}_x = \overrightarrow{\text{constant!}}$$

How can its derivative be  $\neq \vec{0}$

# Second motivation: another issue with polar coordinates

For instance: vector field  $\vec{c}(r, \theta) = c^r(r, \theta)\vec{e}_r(r, \theta) + c^\theta(r, \theta)\vec{e}_\theta(r, \theta)$

To differentiate w.r.t.  $r$  or  $\theta$ , one must not forget to differentiate the basis vectors...

Every time? But that is time-consuming!

There is a nice trick!

# Covariant derivative

Consider the (position dependent:  $P$ ) basis vectors  $\{\vec{e}_i(P)\}$  associated with coordinates  $\{x^i\}$ .

The partial derivative  $\frac{\partial \vec{e}_i(P)}{\partial x^j}$  is itself a vector at  $P$ :

☞ can be written as linear combination of the  $\{\vec{e}_i(P)\}$ :

$$\frac{\partial \vec{e}_i(P)}{\partial x^j} = \Gamma_{ij}^k(P) \vec{e}_k(P)$$

with coefficients  $\Gamma_{ij}^k$ : **Christoffel symbols.** (symmetric under  $i \leftrightarrow j$ )

Consider now a vector field  $\vec{c}(P) = c^i(P) \vec{e}_i(P)$

One has  $\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{dc^i(P)}{dx^j} \vec{e}_i(P)$  with

$$\frac{dc^i(P)}{dx^j} \equiv \frac{\partial c^i(P)}{\partial x^j} + \Gamma_{jk}^i(P) c^k(P)$$

# Covariant derivative

For a vector field  $\vec{c}(P) = c^i(P)\vec{e}_i(P)$  one has  $\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{dc^i(P)}{dx^j}\vec{e}_i(P)$  with the covariant derivatives

$$\frac{dc^i(P)}{dx^j} \equiv \frac{\partial c^i(P)}{\partial x^j} + \Gamma_{jk}^i(P)c^k(P)$$

Proof:

$$\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{\partial c^i(P)}{\partial x^j}\vec{e}_i(P) + c^i(P)\frac{\partial \vec{e}_i(P)}{\partial x^j} = \frac{\partial c^i(P)}{\partial x^j}\vec{e}_i(P) + \underbrace{c^i(P)\Gamma_{ij}^k(P)\vec{e}_k(P)}_{c^i\Gamma_{ij}^k\vec{e}_k = c^k\Gamma_{jk}^i\vec{e}_i}$$

$$\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{\partial c^i(P)}{\partial x^j}\vec{e}_i(P) + c^k(P)\Gamma_{jk}^i(P)\vec{e}_i(P) = \frac{dc^i(P)}{dx^j}\vec{e}_i(P)$$

# Covariant derivative: polar coordinates

From  $\vec{e}_r(r, \theta) = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$  ,  $\vec{e}_\theta(r, \theta) = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$

one computes

$$\frac{\partial \vec{e}_r(r, \theta)}{\partial r} = \vec{0} , \quad \frac{\partial \vec{e}_r(r, \theta)}{\partial \theta} = \frac{1}{r} \vec{e}_\theta(r, \theta) , \quad \frac{\partial \vec{e}_\theta(r, \theta)}{\partial r} = \frac{1}{r} \vec{e}_\theta(r, \theta) , \quad \frac{\partial \vec{e}_\theta(r, \theta)}{\partial \theta} = -r \vec{e}_r(r, \theta)$$

These derivatives are linear combinations  $\frac{\partial \vec{e}_i(P)}{\partial x^j} = \Gamma_{ij}^k(P) \vec{e}_k(P)$  with the Christoffel symbols

$$\Gamma_{rr}^r = \Gamma_{rr}^\theta = 0, \quad \Gamma_{r\theta}^\theta = \Gamma_{\theta r}^\theta = \frac{1}{r}, \quad \Gamma_{\theta\theta}^r = -r, \quad \Gamma_{r\theta}^r = \Gamma_{\theta r}^r = 0, \quad \Gamma_{\theta\theta}^\theta = 0$$

$$\left\{ \begin{array}{l} \frac{dc^i}{dr} = \frac{\partial c^i}{\partial r} + \Gamma_{rk}^i c^k = \frac{\partial c^i}{\partial r} + \Gamma_{r\theta}^i c^\theta = \frac{\partial c^i}{\partial r} + \frac{c^\theta}{r} \delta^{i\theta} \\ \frac{dc^i}{d\theta} = \frac{\partial c^i}{\partial \theta} + \Gamma_{\theta k}^i c^k = \frac{\partial c^i}{\partial \theta} + \Gamma_{\theta r}^i c^r + \Gamma_{\theta\theta}^i c^\theta = \frac{\partial c^i}{\partial \theta} + \frac{c^r}{r} \delta^{i\theta} - r c^\theta \delta^{ir} \end{array} \right.$$

# Covariant derivative: polar coordinates

Consider again the example

$$\vec{c}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta \quad (= \vec{e}_x)$$

i.e. with components  $c^r = \cos \theta$ ,  $c^\theta = -\frac{\sin \theta}{r}$

and let us compute the covariant derivatives w.r.t  $r$  and  $\theta$ :

● w.r.t  $r$ :  $\frac{dc^i}{dr} = \frac{\partial c^i}{\partial r} + \frac{c^\theta}{r} \delta^{i\theta}$  gives  $\frac{dc^r}{dr} = \frac{\partial c^r}{\partial r} = 0$ ,  $\frac{dc^\theta}{dr} = \frac{\partial c^\theta}{\partial r} + \frac{c^\theta}{r} = 0$

● w.r.t  $\theta$ :  $\frac{dc^i}{d\theta} = \frac{\partial c^i}{\partial \theta} + \frac{c^r}{r} \delta^{i\theta} - c^\theta \delta^{ir}$  gives

$$\frac{dc^r}{d\theta} = \frac{\partial c^r}{\partial \theta} - r c^\theta = 0 \quad \text{and} \quad \frac{dc^\theta}{d\theta} = \frac{\partial c^\theta}{\partial \theta} + \frac{c^r}{r} = 0$$

All covariant derivatives vanish!



# Covariant derivative: rules for arbitrary tensor fields

● Scalar field:  $\frac{dc(P)}{dx^j} \equiv \frac{\partial c(P)}{\partial x^j}$

● "Contravariant vector":  $\frac{dc^i(P)}{dx^j} \equiv \frac{\partial c^i(P)}{\partial x^j} + \Gamma_{jk}^i(P)c^k(P)$

● "Covariant vector":  $\frac{dc_i(P)}{dx^j} \equiv \frac{\partial c_i(P)}{\partial x^j} - \Gamma_{ij}^k(P)c_k(P)$

● Arbitrary (m contra-, n covariant indices) tensor field:

$$\begin{aligned} \frac{d\mathbf{T}_{j_1 \dots j_n}^{i_1 \dots i_m}(P)}{dx^k} &= \frac{\partial \mathbf{T}_{j_1 \dots j_n}^{i_1 \dots i_m}(P)}{\partial x^k} + \Gamma_{kl}^{i_1}(P)\mathbf{T}_{j_1 \dots j_n}^{li_2 \dots i_m}(P) + \dots + \Gamma_{kl}^{i_m}(P)\mathbf{T}_{j_1 \dots j_n}^{i_1 \dots i_{m-1}l}(P) \\ &\quad - \Gamma_{j_1 k}^l(P)\mathbf{T}_{lj_2 \dots j_n}^{i_1 \dots i_m}(P) - \dots - \Gamma_{j_n k}^l(P)\mathbf{T}_{j_1 \dots j_{n-1}l}^{i_1 \dots i_m}(P) \end{aligned}$$

# A last formula

To compute the Christoffel symbols with the help of the metric tensor and its inverse 🖱️ no need to go back to the expression of the basis vectors  $\{\vec{e}_i\}$  in terms of those of the Cartesian / Minkowski basis.

$$\Gamma_{ij}^k(P) = \frac{1}{2} g^{kl}(P) \left[ \frac{\partial g_{jl}(P)}{\partial x^i} + \frac{\partial g_{il}(P)}{\partial x^j} - \frac{\partial g_{ij}(P)}{\partial x^l} \right]$$

🖱️ easy to automatize (symbolic algebra programs)