

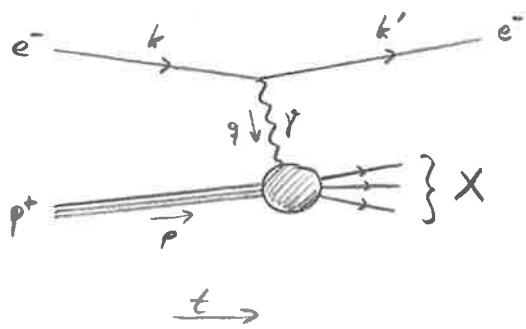
## 5. Deep inelastic scattering (DIS)

We now (temporarily) go back to tree-level phenomenology.

Q.: are quarks real physical constituents of hadrons,  
or just a mathematical convenience for describing the hadron's wavefct?

→ DIS gives information on internal proton structure

→ structure functions, Bjorken scaling, parton distribution fcts



- photon virtuality -  $Q^2 = -q^2$  ✓ ( $Q^2 \sim \frac{1}{\lambda^2}$ )  
controls resolving power of photon

$Q^2 \ll \frac{1}{R_p^2} \rightarrow$  elastic ep scattering  
proton radius

$Q^2 \gg \frac{1}{R_p^2} \rightarrow$  resolve p constituents,  
elastic e-g scattering

we are interested in deep ( $Q^2 \gg M_p^2$ ) inelastic ( $(p_{\gamma p})^2 \gg D_p^2$ ) scattering

### 5.1 Structure functions

- want to describe process with Lorentz-invariant variables

$$\text{center-of-mass energy } s \equiv (k+p)^2$$

at fixed  $s$ , scattered  $e^-$  has 2 non-trivial variables ( $E, \theta$ );

$$\text{use } Q^2 = -q^2, \quad x \equiv \frac{Q^2}{2p \cdot q}$$

$$(\text{other choices: } W^2 \equiv (p_{\gamma p})^2 = Q^2 \frac{1-x}{x} \quad (\text{invariant mass of } \gamma p\text{-system})$$

$$y \equiv \frac{p \cdot q}{p \cdot k} = \frac{Q^2}{xs} \quad \parallel$$

$$\text{kinematic limits: } Q^2 < s, \quad x > \frac{Q^2}{s}$$

$$(\text{of course } k^2 = m_e^2, \quad p^2 = M_p^2; \text{ for } Q^2 \gtrsim D_p^2, \text{ also } Z^0 \text{ exchange})$$

- know nothing about detailed structure of proton

→ parametrize  $M = \frac{e^4}{Q^4} \text{tr}(k' g^{\mu\nu} k' g^{\nu\mu}) T_\mu(p, q; \{p_x\}) T_\nu^*(p, q; \{p_x\})$

$$\Rightarrow \frac{1}{4} |M|^2 = \frac{1}{4} \frac{e^4}{Q^4} \underbrace{\text{tr}(k' g^{\mu\nu} k' g^{\nu\mu})}_{= L^{\mu\nu}} T_\mu(p, q; \{p_x\}) T_\nu^*(p, q; \{p_x\}) \\ = L^{\mu\nu} = 4(k'^\mu k'^\nu + k^\nu k'^\mu - k \cdot k' g^{\mu\nu}) \quad (\text{see p. 47})$$

in complete analogy to  $e^- e^- \rightarrow \mu^+ \mu^-, q\bar{q}$  etc.

- for total cross section, need to integrate over phase space

$$\int d\Phi_{X+1} = \underbrace{\int dQ^2 dx \frac{Q^2}{16\pi^2 s x^2}}_{\text{electron kinematics}} \underbrace{\int d\Omega_X}_{n\text{-body phase space for } X}$$

for an inclusive process (don't measure  $X \rightarrow$  sum over elem.),

$$\frac{d^2\sigma}{dQ^2 dx} = \frac{1}{2s} \frac{Q^2}{16\pi^2 s x^2} \underbrace{\sum_X \int d\Omega_X \frac{1}{4} |M|^2}_{= \frac{1}{4} \frac{e^4}{Q^4} L^{\mu\nu} \underbrace{\sum_X \int d\Omega_X T_\mu(p, q; \{p_x\}) T_\nu^*(p, q; \{p_x\})}_{= H_{\mu\nu}},}$$

- consider  $H_{\mu\nu}$ : have summed and integrated all  $X$  dependence, so  $H_{\mu\nu}(p, q)$ , must be symm. in  $\mu, \nu$  ((parity cons. in QED, a.c.d))

$$\Rightarrow H_{\mu\nu} = -H_1 \delta_{\mu\nu} + H_2 \frac{p_\mu p_\nu}{Q^2} + H_3 \frac{q_\mu q_\nu}{Q^2} + H_4 \frac{p_\mu q_\nu + q_\mu p_\nu}{Q^2}$$

where  $H_i$  are scalar fits,  $H_i(Q^2 = -Q^2, p \cdot q = \frac{Q^2}{2x}, \rho^2 = \frac{Q^2}{s})$   
neglect in DIS

((including  $Z^0$  exchange, ... +  $H_5 \epsilon_{\mu\nu\rho\sigma} \frac{p_\rho q_\sigma}{Q^2}$  ))

$$\Rightarrow L^{\mu\nu} H_{\mu\nu} = 8(k') H_1 + 8 \frac{(\rho t)(\rho t')}{Q^2} H_2 \stackrel{\text{(check ??!)}}{\downarrow} \downarrow H_3 + 0 \cdot H_4$$

used  $d=4$ , neglected  $\rho^2$

$$\text{now } Q^2 = -q^2 = -(t-t')^2 = 2tt' - 2q^2$$

$$s = (\rho+t)^2 = 2\rho t + \cancel{2t^2}$$

$$\rho t' = \rho(t-t) = \rho t(1 - \frac{t}{\rho t}) = \rho t(1-q)$$

$$= 4Q^2 H_1 + 2 \frac{s^2}{Q^2} (1-q) H_2$$

- $$\frac{d^2\sigma}{dQ^2 dx} = \frac{1}{25} \frac{Q^2}{16\pi^2 x^2} \frac{1}{4} \frac{\alpha^2/16\pi^2}{Q^4} \left( 4Q^2 H_1 + 2 \frac{\alpha^2}{Q^2} (1-y) H_2 \right); \text{ def } \begin{array}{l} H_1 = 8\pi F_1 \\ H_2 = 16\pi x F_2 \end{array}$$

$$= \frac{4\pi\alpha^2}{x Q^4} \left\{ x y^2 F_1(x, Q^2) + (1-y) F_2(x, Q^2) \right\}$$

→ amazing: w/o knowing ep interaction, derived s-dependence of  $\sigma$ ! (recall  $q = Q^2/x_S$ )

→ the  $F$ 's are called "structure functions" of the proton

sometimes, see  $F_T(x, Q^2) = 2x F_1(x, Q^2)$  "transverse"  
 $F_L(x, Q^2) = F_2(x, Q^2) - 2x F_1(x, Q^2)$  "longitudinal"

so 
$$\frac{d^2\sigma}{dQ^2 dx} = \frac{2\pi\alpha^2}{x Q^4} \left\{ (1+(1-y)^2) F_T(x, Q^2) + 2(1-y) F_L(x, Q^2) \right\}$$
  

$$= \frac{2\pi\alpha^2}{x Q^4} \left\{ (1+(1-y)^2) F_2(x, Q^2) - y^2 F_L(x, Q^2) \right\}$$

useful since (for most current data)  $y^2 \ll 1$

- have isolated all non-trivial  $x, Q^2$  dependence into  $F_2, F_L$  ;  
 but still don't know anything about these fcts.

Assumption: interaction of  $g$  with innards of proton does not involve any dimensionful scale

→ dim-less  $F$ 's can not depend on dimensionful parameter  $Q^2$

$$\Rightarrow \frac{d^2\sigma}{dQ^2 dx} = \frac{2\pi\alpha^2}{x Q^4} \left\{ (1+(1-y)^2) F_2(x) - y^2 F_L(x) \right\} \quad \text{Bjorken scaling}$$

→ experimentally, this is true (to a pretty good approximation);  
 but proton consists of quarks, bound at distance scale  $\sim 1/\mu_p$ ,  
 so how can the interaction possibly be  $\mu_p$ -indep. ?!

→ answer via parton model