

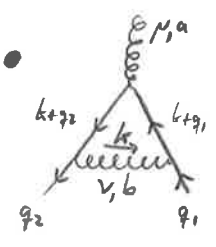
Fy par,  $\frac{1}{(x)} = \int_0^1 dx \frac{1}{[(1-x)l^2 + x(l^2 + 2lq_1 + q_1^2 - m^2)]^2} = \int_0^1 dx \frac{1}{[(l+xq_1)^2 + x(1-x)q_1^2 - xm^2]^2}$

shift  $l \rightarrow l - xq_1$   
 $-ig^2 \frac{N_c^2 - 1}{2N_c} \int \frac{d^d k}{(2\pi)^d} \int_0^1 dx \frac{(2-d)(k + (1-x)q_1) + md}{(l^2 - \Delta)^2}$

Wick rotate, use basic tadpole from pg 31

$= -ig^2 \frac{N_c^2 - 1}{2N_c} \int_0^1 dx \underbrace{I_2^0(xm^2 - x(1-x)q_1^2)}_{\approx \frac{1}{\epsilon} \frac{1}{(4m)^2}, \text{ see pg. 34}} \{ (2-d)(1-x)q_1 + md \}$

$\approx i \frac{g^2}{16\pi^2} \frac{N_c^2 - 1}{2N_c} \frac{1}{\epsilon} \left\{ \frac{1}{\xi} - \frac{1}{4m} + O(\epsilon) \right\}$   
 for general value of  $\xi$



$\int \frac{d^d k}{(2\pi)^d} (ig)^3 T^b T^a T^b \frac{\gamma^\nu i(k + \cancel{q}_2 + m) \gamma^\mu i(k + \cancel{q}_1 + m) \gamma_\nu (-i)}{((l+q_2)^2 - m^2) ((l+q_1)^2 - m^2) l^2}$

$\gamma^{\mu\nu\sigma} = (2g^{\mu\nu} - \gamma^\mu \gamma^\nu) \gamma^\sigma = 2\gamma^{\sigma\nu} - \gamma^\nu \gamma^{\mu\sigma} = 4g^{\nu\sigma} - (4-d)\gamma^{\nu\sigma}$   
 $= 2\gamma^{\sigma\nu} - \gamma^{\nu\sigma} = (2-d)\gamma^\sigma$

$\gamma^{\mu\nu\sigma\sigma} = (2g^{\mu\nu} - \gamma^\mu \gamma^\nu) \gamma^{\sigma\sigma} = 2\gamma^{\sigma\sigma} \gamma^\nu - \gamma^\nu (4g^{\sigma\sigma} - (4-d)\gamma^{\sigma\sigma}) = -2\gamma^{\sigma\sigma\nu} + (4-d)\gamma^{\nu\sigma\sigma}$   
 $= 2\gamma^{\sigma\sigma} - \gamma^{\sigma\sigma}$

$T_{ij}^b T_{jb}^a T_{ai}^b = T_{ja}^a \frac{1}{2} (\delta_{il} \delta_{jb} - \frac{1}{N_c} \delta_{ij} \delta_{ba}) = \frac{1}{2} \delta_{il} T_{ji}^a - \frac{1}{2N_c} T_{il}^a = -\frac{1}{2N_c} T^a$

$g^3 (-\frac{1}{2N_c} T^a) \int \frac{d^d k}{(2\pi)^d} \frac{N^\mu}{((l+q_2)^2 - m^2) ((l+q_1)^2 - m^2) l^2}$

$N^\mu = m^2(2-d)\gamma^\mu + m(l+q_2)_\nu (4g^{\nu\mu} - (4-d)\gamma^{\nu\mu}) + m(l+q_1)_\nu (4g^{\mu\nu} - (4-d)\gamma^{\mu\nu}) + (l+q_2)_\nu (l+q_1)_\sigma [-2\gamma^{\sigma\nu\mu} + (4-d)\gamma^{\nu\mu\sigma}]$

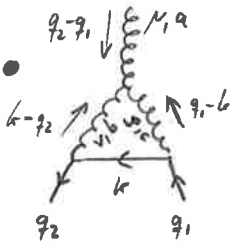
superficial degree of divergence of this integral:  $\frac{6}{d} \rightarrow \log!$

$\Rightarrow$  look at  $l \gg \{q_1, q_2, m\}$  only, to extract this leading log div

$\sim g^3 (-\frac{1}{2N_c} T^a) \int \frac{d^d k}{(2\pi)^d} \frac{[-2\gamma^{\sigma\nu\mu} + (4-d)\gamma^{\nu\mu\sigma}] l_\nu l_\sigma}{(l^2)^3} + \text{const.}$

numerator  $\rightarrow [\dots] \frac{g_{\nu\sigma} l^2}{d} = \frac{l^2}{d} [-2\gamma^{\sigma\mu}{}_\sigma + (4-d)\gamma^{\nu\mu}{}_\nu] \stackrel{\text{pg. 34}}{=} \frac{l^2}{d} (2-d)^2 \gamma^{\mu\nu}$

$g^3 (-\frac{1}{2N_c} T^a) \frac{(2-d)^2}{d} \gamma^{\mu\nu} \int \frac{d^d k}{(2\pi)^d} \frac{1}{(l^2)^2} = i I_2^0(0)$  (after Wick rot.)



$$\int \frac{d^d k}{(2\pi)^d} (ig)^2 T^b T^c \frac{\gamma_\nu i(\not{k} + m) \gamma_\sigma (-i) (-i)}{(k^2 - m^2) (k - q_1)^2 (k - q_2)^2} g f^{abc} \epsilon$$

$$* \left[ (2q_2 - q_1 - k)^\sigma g^{\mu\nu} + (2k - q_1 - q_2)^\mu g^{\nu\sigma} + (2q_1 - q_2 - k)^\nu g^{\sigma\mu} \right]$$

$$\Gamma T^b T^c f^{abc} = \frac{1}{2} [T^b, T^c] f^{abc} = \frac{i}{2} f^{bcd} T^d f^{abc} = \frac{i}{2} N_c T^a$$

=  $N_c \delta^{ad}$  (see pg. 33)

again, extract only leading log div,  $k \gg \{q_1, q_2, m\}$

$$\sim -g^3 \frac{N_c}{2} T^a \int \frac{d^d k}{(2\pi)^d} \frac{\gamma_\nu \not{k} \gamma_\sigma [-k^\sigma g^{\mu\nu} + 2k^\mu g^{\nu\sigma} - k^\nu g^{\sigma\mu}]}{(k^2)^3} + \text{const.}$$

replace  $k^\mu k^\nu \rightarrow g^{\mu\nu} \frac{k^2}{d}$

numerator  $\gamma_\nu \not{k} \gamma_\sigma [-k^\sigma g^{\mu\nu} + 2k^\mu g^{\nu\sigma} - k^\nu g^{\sigma\mu}]$

$$= -\gamma^\mu \not{k} \gamma_\sigma + 2\gamma^{\sigma\mu} \not{k} - \gamma^\nu \not{k} \gamma_\sigma = -d\gamma^\mu + 2(2-d)\gamma^\mu - d\gamma^\mu = 4(1-d)\gamma^\mu$$

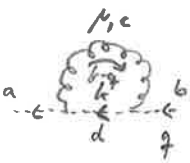
$$= -g^3 2N_c T^a \frac{1-d}{d} \gamma^\mu \int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2)^2} = i I_2^0(0) \quad (\text{after Wick rot.})$$

• sum of last two diagrams

$$\text{triangle} + \text{triangle} = -ig^3 T^a \gamma^\mu I_2^0(0) \left\{ \frac{1}{2N} \frac{(2-d)^2}{d} + 2N \frac{1-d}{d} \right\} + \text{const.}$$

$$d=4-2\epsilon \rightsquigarrow -ig T^a \gamma^\mu \frac{g^2}{16\pi^2} \frac{1}{\epsilon} \left\{ \frac{1}{2N} - \frac{3}{4} [2N + \mathcal{O}(\epsilon)] \right\}$$

$\rightarrow (\xi+1)$  for general value of  $\xi$



$$= \int \frac{d^d k}{(2\pi)^d} (-g f^{acd} \gamma^\mu) (-g f^{deb} \gamma_\nu) \frac{i}{k^2} \frac{-i}{(k-q)^2}$$

$$= -g^2 f^{acd} f^{deb} \int \frac{d^d k}{(2\pi)^d} \frac{\gamma^\mu \not{k} \gamma_\nu}{k^2 (k-q)^2}$$

=  $N_c \delta^{ab}$  (pg. 33)

denominator invariant under  $k \rightarrow q-k$

write numerator  $k = \frac{1}{2}k + \frac{1}{2}k \rightarrow \frac{1}{2}k + \frac{1}{2}(q-k) = \frac{1}{2}q$

$$= -g^2 N_c \delta^{ab} \frac{g^2}{2} \int \frac{d^d k}{(2\pi)^d} \frac{1}{k^2 (k-q)^2}$$

extract leading log div,  $k \gg q$ ; Wick rot; use tadpole

$$\sim -ig^2 N_c \delta^{ab} \frac{g^2}{2} I_2^0(0) + \text{const.}$$

$$d=4-2\epsilon \rightsquigarrow -i \frac{g^2}{16\pi^2} \delta^{ab} g^2 \frac{1}{\epsilon} \left\{ [2] \frac{N_c}{4} + \mathcal{O}(\epsilon) \right\}$$

$\rightarrow (3-\xi)$  for general value of  $\xi$

### 3.3 one-loop counterterms in QCD

→ now, renormalize the theory.

use freedom of redefining fields, parameters/couplings

schematically:  $\phi_B = \sqrt{Z_\phi} \phi_R$ ,  $\phi \in \{\psi, A, c\}$

$\lambda_B = Z_\lambda \lambda_R$ ,  $\lambda \in \{m, g, \xi\}$

B = "bare"

R = "renormalized"

where the multiplicative renormalization factors  $Z_i$

depend on the renormalized parameters (and the dimension  $d$ ),

and are taken to be dimensionless,  $Z_i = 1 + \delta Z_i$ ,  $\delta Z_i \sim g^2$  (see below)

• recall (pg. 28)

$$\mathcal{L}_B = \bar{\psi}_B (i \gamma^\mu D_\mu - m_B) \psi_B - \frac{1}{4} F_{\mu\nu}^i F^{\mu\nu i} - \frac{1}{2\xi} (\partial^\mu A_\mu^a)^2 + \bar{c}^a (-\partial^\mu D_\mu^{ac}) c^a$$

$\int_{\psi} \int_{A} - i g A_\mu^a T^a$       $\int_{A} \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c$       $\int_{c} \partial_\mu^a c^a + g f^{abc} A_\mu^b c^c$

(in this line, all  $\psi, A, c, m, g, \xi$  should have an index B)

$$\sqrt{Z_\psi} \bar{\psi} i \not{\partial} \psi - \sqrt{\frac{Z_m Z_\psi}{m}} m \bar{\psi} \psi + \sqrt{\frac{Z_g Z_\psi Z_A^2}{g}} g \bar{\psi} \gamma^\mu T^a \psi A_\mu^a$$

$$-\sqrt{\frac{Z_A}{4}} (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a)^2 - \sqrt{\frac{Z_\xi Z_A^{-1}}{2\xi}} (\partial^\mu A_\mu^a)^2 - \sqrt{\frac{Z_g Z_A^{3/2}}{2}} g f^{abc} A_\mu^b A_\nu^c (\partial^\mu A^\nu - \partial^\nu A^\mu)$$

$$-\sqrt{\frac{Z_g^2 Z_A^2}{4}} g^2 f^{abc} A_\mu^b A_\nu^c f^{ade} A^\mu A^\nu - \sqrt{\frac{Z_c}{\xi}} \bar{c}^a \partial^\mu c^a - \sqrt{\frac{Z_g Z_c Z_A^{3/2}}{g}} g f^{abc} \bar{c}^a \partial^\mu c^b A_\mu^c$$

(here, without  $Z$ 's: index B; with  $Z$ 's: index R for all  $\psi, A, c, m, g, \xi$ )

$\mathcal{L}_R + \mathcal{L}_{c.t.}$  ← counterterms

$$= (Z_\psi - 1) \bar{\psi} (i \not{\partial} - \frac{Z_m Z_\psi - 1}{Z_\psi - 1} m) \psi + (Z_g Z_\psi Z_A^2 - 1) \text{cancel}$$

$$- (Z_A - 1) \left[ \frac{1}{4} (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a)^2 + \frac{Z_A Z_\xi^{-1} - 1}{Z_A - 1} \frac{1}{2\xi} (\partial^\mu A_\mu^a)^2 \right] + (Z_g Z_A^{3/2} - 1) \text{cancel}$$

$$+ (Z_g^2 Z_A^2 - 1) \text{cancel} - (Z_c - 1) \bar{c}^a \partial^\mu c^a + (Z_g Z_c Z_A^{3/2} - 1) \text{cancel}$$

(now, all indices are R, and omitted)