

have used the following expansions:

(*) $y^{-\epsilon} = 1 - \epsilon \ln y + \mathcal{O}(\epsilon^2)$
 $\frac{1}{\epsilon} B\left(\frac{1-\epsilon}{2}, \frac{\epsilon}{2}\right) = \frac{1}{\epsilon} + \ln 2 + \mathcal{O}(\epsilon)$

(**) $\frac{1}{\epsilon} B\left(\frac{1+\epsilon}{2}, \frac{1+\epsilon}{2}\right) = \frac{\pi}{\epsilon} + \mathcal{O}(\epsilon^2)$
 $\frac{1}{(2\pi)^{\epsilon}} \int (1+\epsilon) = \frac{1}{\epsilon} + \gamma - \ln(2\pi) + \mathcal{O}(\epsilon)$

c) get $n=3, n=5$ -cases via integration

$$h_3(z) = \frac{\pi^2}{2} \left(\frac{1}{6} - z - z^2 \left(\ln \frac{z}{2} + \gamma - \frac{1}{2} \right) - \sum_{j=1}^{\infty} \frac{c_j}{j!} \int (2j-1) z^{2j+2} \right) \quad (76)$$

$$h_5(z) = \frac{\pi^4}{48} \left(\frac{1}{15} - z^2 + 4z^3 + 3z^4 \left(\ln \frac{z}{2} + \gamma - \frac{3}{4} \right) + \sum_{j=1}^{\infty} \frac{6c_j}{(j+1)j!} \int (2j-1) z^{2j+4} \right) \quad (77)$$

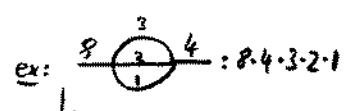
etc...

• 2nd order of pert. expansion (NNLO)

$$Z_{int} = 1 + 3 \mathcal{O} + \frac{1}{2} (9 \mathcal{O}^2 + 72 \mathcal{O} \mathcal{O} + 24 \mathcal{O}^3) + \mathcal{O}(\lambda^3)$$

$$\ln Z_{int} = 3 \mathcal{O} + (36 \mathcal{O}^2 + 12 \mathcal{O} \mathcal{O}) + \mathcal{O}(\lambda^3) \quad (78)$$

⇒ products of vacuum diag's cancelled!



$$\tilde{G}_2 = \frac{1}{Z_{int}} \left(- + (12 \mathcal{O} + 3 \mathcal{O}^2) + \frac{1}{2} (288 \mathcal{O} + 288 \mathcal{O} \mathcal{O} + 192 \mathcal{O}^2 + 72 \mathcal{O}^3 + 9 \mathcal{O}^4 + 72 \mathcal{O} \mathcal{O}^2 + 24 \mathcal{O}^3) + \mathcal{O}(\lambda^3) \right)$$

$$= - + 12 \mathcal{O} + (144 \mathcal{O} + 144 \mathcal{O} \mathcal{O} + 96 \mathcal{O}^2) + \mathcal{O}(\lambda^3) \quad (79)$$

⇒ vacuum subdiagrams cancelled!

these cancellations are a general feature of the perturbative expansion.

⇒ present a general (all-order) proof! → FT²₃₆, FT²₃₇

($\ln(1+x) = x - \frac{x^2}{2} + \dots$)

($\frac{1}{1+x} = 1 - x + x^2 - \dots$)

proof of vac-subbing-cancellation

a) Z_{nt}

$$Z_{nt} = [e^{Q \cdot U_0}]_{j=0} = 1 + \sum_{n=1}^{\infty} \frac{1}{n!} [a_1 \dots a_n U_0]_{j=0}$$

ground structure: $\forall Q$ (from $S^{(j)}$ part) contain 1-op's.

these can a) pull a j -th out of $U_0 = e^{\sum_{i,j} \delta_{ij} \cdot a_i}$
 b) left a j denominator

in the end, no j 's or l 's left, $U_0[0]=1$
 \rightarrow values count in all possible ways.

$$= 1 + \sum_{n=1}^{\infty} \frac{1}{n!} \sum_{m=1}^n \frac{1}{m!} \left[\sum_{\text{partitions}} [Q^{r_1} U_0]_{j=0} \dots [Q^{r_m} U_0]_{j=0} \right] \equiv V_{n,m}$$

\uparrow
 # groups \rightarrow to partitions, enumerated Q 's over no numbered groups
 in unnumbered groups overcount!


$$= 1 + \sum_{n=1}^{\infty} \frac{1}{n!} \sum_{m=1}^n \frac{1}{m!} V_{n,m}$$

$$V_{n,1} = [Q^n U_0]_{j=0}$$

$$V_{n,n} = n! ([Q U_0]_{j=0})^n$$

$$V_{n,m \geq 2} = \sum_{r=1}^{n-(m-1)} \binom{n}{r} [Q^r U_0]_{j=0} V_{n-r, m-1} \quad (\text{recursion})$$

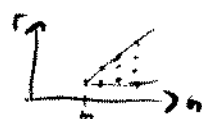
\uparrow
 # of ways to pick r out of n Q 's
 (and put in additional group)



$$\rightarrow \sum_{n=1}^{\infty} \sum_{m=1}^n = \sum_{m=1}^{\infty} \sum_{n=m}^{\infty}$$

$$= 1 + \sum_{n=1}^{\infty} \frac{1}{n!} \left[\sum_{m=1}^n \frac{1}{m!} V_{n,m} \right] S_m, \quad S_1 = \sum_{n=1}^{\infty} \frac{1}{n!} V_{n,1} = \sum_{n=1}^{\infty} \frac{1}{n!} [Q^n U_0]_{j=0} = [(e-1)U_0]_{j=0}$$

$$= \sum_{n=2}^{\infty} \frac{1}{n!} \sum_{r=1}^{n-(m-1)} \frac{n!}{r!(n-r)!} [Q^r U_0]_{j=0} V_{n-r, m-1}$$



$$\rightarrow \sum_{n=m}^{\infty} \sum_{r=1}^{n-(m-1)} = \sum_{r=1}^{\infty} \sum_{n=r+(m-1)}^{\infty}$$

$$= \sum_{r=1}^{\infty} \frac{1}{r!} [Q^r U_0]_{j=0} \sum_{n=r+(m-1)}^{\infty} \frac{1}{(n-r)!} V_{n-r, m-1} = \sum_{n=m-1}^{\infty} \frac{1}{n!} V_{n, m-1} = S_{m-1}$$

$$S_m = S_1 \cdot S_{m-1} \rightsquigarrow S_m = S_1^m$$

$$Z_{int} = 1 + \sum_{m=1}^{\infty} \frac{1}{m!} S_1^m = e^{S_1}$$

$$\Rightarrow \ln Z_{int} = S_1 = [(e^{\alpha}-1)W_0]_{con} \quad (80)$$

\rightsquigarrow there are no products of clouds in $\ln Z_{int}$!!!

((does this only apply to a one-index-theory ??))

6) $W[j]$

$$W[j] = \frac{1}{Z_{int}} e^{\alpha W_0[j]} \quad (\text{not } j=0), W[0]=1$$

$$= \frac{1}{Z_{int}} \left(W_0[j] + \sum_{n=1}^{\infty} \frac{1}{n!} [Q_1 \dots Q_n W_0[j]] \right)$$

philo: p values are not part of vac-subdiag's (=clouds), the rest ($n-p$) connects to clouds

$$\sum_{p=0}^{n-1} (Q^p W_0)_{\text{cloud-free}} \binom{n}{p} \sum_{m=1}^{n-p} \frac{1}{m!} \sum_{\text{conn}} [Q^m W_0]_{con} - [Q^n W_0]_{con}$$

$\sum_i c_i = n-p, c_i \geq 1$

$$= \frac{1}{Z_{int}} \left((e^{\alpha W_0})_{cf} + \sum_{n=1}^{\infty} \frac{1}{n!} \sum_{p=0}^{n-1} (Q^p W_0)_{\text{cloud-free}} \sum_{m=1}^{n-p} \frac{1}{m!} V_{n-p,m} \right)$$

$$\sum_{p=0}^{n-1} \sum_{m=1}^{n-p} \sum_{n=1}^{\infty} \sum_{m=1}^{n-p} \frac{1}{n!} \frac{1}{m!} V_{n-p,m}$$

$$= \frac{1}{Z_{int}} \left((e^{\alpha W_0})_{cf} + \sum_{p=0}^{\infty} \frac{1}{p!} (Q^p W_0)_{cf} \sum_{m=1}^{\infty} \frac{1}{m!} \sum_{n=p+m}^{\infty} \frac{1}{(n-p)!} V_{n-p,m} \right)$$

$$= \frac{1}{Z_{int}} \left((e^{\alpha W_0})_{cf} + (e^{\alpha W_0})_{cf} (e^{S_1} - 1) \right) \quad S_1 = [(e^{\alpha}-1)W_0]_{con}$$

$$= \frac{e^{S_1}}{Z_{int}} (e^{\alpha W_0})_{cf}$$

$$W[j] = (e^{\alpha W_0})_{\text{cloud-free}} \quad (81)$$

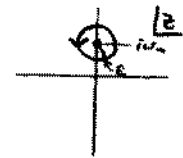
one can do more:
 • $\ln W$ generates labelled con
 (eg: $\begin{matrix} \text{not con} \\ \text{not con} \end{matrix} \quad \begin{matrix} \text{labelled con} \\ \text{labelled con} \end{matrix}$)

(one more) technical tool:

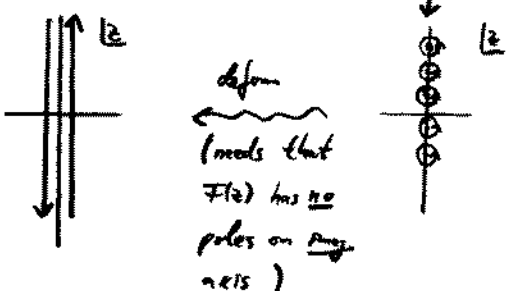
• frequency sums by contour integration

(have to handle, eg., $T \sum_{n=-\infty}^{\infty} \frac{1}{k^2}$, $k = (i\omega_n, \vec{k})$, $\omega_n = \pi(2n + \frac{1}{2})T$)
 $k^2 = (i\omega_n)^2 - \vec{k}^2$ (Fermi)

more generally, $T \sum_{n=-\infty}^{\infty} \frac{F(i\omega_n)}{k}$ analytic fct (no poles on imag. axis, see below)
 can trade \sum_n for integral \rightarrow easier?!

1. $F(i\omega_n) = \frac{1}{2\pi i} \oint_n dz \frac{F(z)}{z - i\omega_n}$ (make harder)  (82)

2. $\frac{T}{z - i\omega_n}$ (play) inf. small! $= \frac{1}{e^{\beta(z - i\omega_n)} - 1} = \frac{1}{\eta e^{\beta z} - 1}$
 $= e^{-i\beta(2n + \frac{1}{2})T} = \pm 1 \rightarrow \eta^2 = 1$
 $= \eta^n(z)$, $n(z) = \frac{1}{e^{\beta z} - \eta}$ (83)

3. $T \sum_{n=-\infty}^{\infty} F(i\omega_n) \stackrel{1.}{=} T \sum_n \frac{1}{2\pi i} \oint_n dz \frac{F(z)}{z - i\omega_n} \stackrel{2.}{=} \frac{T}{2\pi i} \sum_n \oint_n dz F(z) n(z)$ (deform contour)

 (needs that $F(z)$ has no poles on imag. axis)

$\Rightarrow T \sum_{n=-\infty}^{\infty} F(i\omega_n) = \frac{T}{2\pi i} \int_{\mathcal{C}} dz F(z) n(z)$ (84)

4. (rewrite) in left integral ($\int_{\mathcal{C}}$), let $z \rightarrow -z$
 $\int_{\mathcal{C}} dz F(z) n(z) = \int_{i\omega_2}^{-i\omega_2} dz F(z) n(z) = - \int_{-i\omega_2}^{i\omega_2} dz F(-z) n(-z) = - \int_{\mathcal{C}} dz F(-z) n(-z)$
 note $n(-z) = \frac{1}{e^{-\beta z} - \eta} = -\eta - \frac{1}{e^{\beta z} - \eta} = -\eta - n(z)$

$\Rightarrow T \sum_{n=-\infty}^{\infty} F(i\omega_n) = \frac{1}{2\pi i} \int_{\mathcal{C}} dz F(-z) + \frac{T}{2\pi i} \int_{\mathcal{C}} dz (F(z) + F(-z)) n(z)$
 $= \frac{1}{2\pi i} \int_{-i\omega}^{i\omega} dz \frac{1}{2} (F(z) + F(-z)) + \frac{T}{2\pi i} \int_{-i\omega}^{i\omega} dz (F(z) + F(-z)) n(z)$
 "T=0" - piece (85)

version of Koba (3.10) \rightarrow