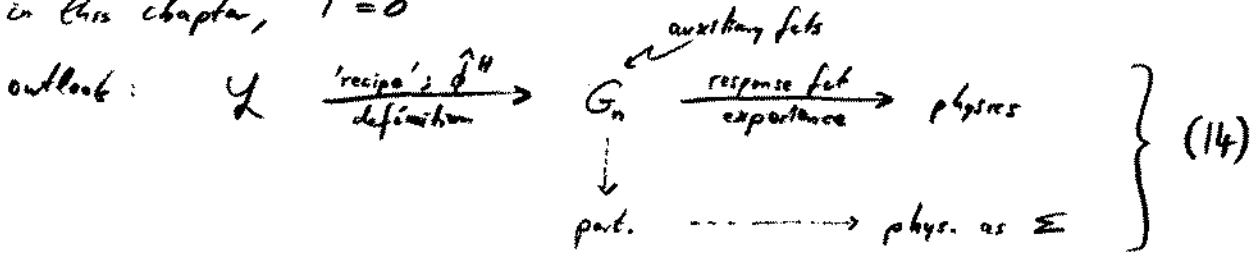


3) (review of) Quantum Field Theory

from now on, $\hbar = c = k_B = 1$

in this chapter, $T=0$



(for a quick intro, take) simplest Field Theory (?)

• free bosonic field in 0+1 D

$$\left. \begin{aligned} \int d^3r \rightarrow 1, \quad \sqrt{m^2 + k^2} \rightarrow m \\ \mathcal{L} = L = \frac{1}{2} \dot{\phi}^2 - \frac{1}{2} m^2 \phi^2 \end{aligned} \right\} (15)$$

(eqn of motion via $\partial_\mu \partial_{\mu\nu} \mathcal{L} = \partial_\nu \mathcal{L} \Rightarrow \ddot{\phi} = -m^2 \phi$)

\rightarrow harm. osc! $\phi \hat{=} q, \quad m \hat{=} \omega, \quad \pi \hat{=} p$

(conj. momenta via $\pi_\mu = \partial_{\dot{\phi}_\mu} \mathcal{L} \Rightarrow \pi = \dot{\phi}$)

$$H_{\text{class.}} = \pi \dot{\phi} - L = \frac{1}{2} \pi^2 + \frac{1}{2} m^2 \phi^2$$

$$\pi \rightarrow -i \partial_\phi, \quad H_{\text{class.}} \rightarrow H = -\frac{1}{2} \partial_\phi^2 + \frac{1}{2} m^2 \phi^2$$

"QM at one point, in 'field direction'"

equal time commutators $[\phi, \pi] = i, \text{ rest } = 0$

($[\phi_S(t, \vec{r}), \pi_S(t', \vec{r}')] = i \delta(\vec{r} - \vec{r}') \delta(t - t') \delta_{SS'}$, $\{\dots\}$ for Fermi)

algebraical method

$$b = \sqrt{\frac{m}{2}} \phi + \frac{1}{\sqrt{2m}} \partial_\phi, \quad b^\dagger = \dots$$

$$\rightarrow \phi = \frac{1}{\sqrt{2m}} (b + b^\dagger) \quad \leftarrow \text{coordinate operator}$$

$$\rightarrow H = \frac{m}{2} (b b^\dagger + b^\dagger b), \quad [b, b^\dagger] = 1$$

Schrödinger functional
 $H\psi(\phi) = E\psi(\phi), \psi_0, E_0$
 \downarrow
 $\psi[\phi(t=0)]$

change of representation ("2nd quant.")

$|n\rangle$ cons (complete orthonormal system), $|n\rangle = \frac{1}{\sqrt{n!}} (\hat{b}^\dagger)^n |0\rangle$, $\hat{b}|0\rangle=0$

$$\begin{aligned} \langle n' | \hat{b} | n \rangle &= \langle n' | \frac{1}{\sqrt{n!}} \hat{b} (\hat{b}^\dagger)^n | 0 \rangle, \quad [\hat{b}, \hat{b}^\dagger] = 1 \\ &= \langle n' | \left(\frac{1}{\sqrt{n!}} (\hat{b}^\dagger)^{n-1} + \frac{1}{\sqrt{n!}} \hat{b}^\dagger \hat{b} (\hat{b}^\dagger)^{n-1} \right) | 0 \rangle \\ &\vdots \\ &= \langle n' | \left(\frac{n}{\sqrt{n!}} (\hat{b}^\dagger)^{n-1} + \frac{1}{\sqrt{n!}} (\hat{b}^\dagger)^n \hat{b} \right) | 0 \rangle \\ &= \sqrt{n} \langle n' | n-1 \rangle \end{aligned}$$

$\Rightarrow \hat{b}|n\rangle = \sqrt{n}|n-1\rangle$ (and $\hat{b}^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle \Rightarrow \hat{b}^\dagger \hat{b}|n\rangle = n|n\rangle$)

$$\begin{aligned} \hat{H} &= \frac{m}{2} (\hat{b} \hat{b}^\dagger + \hat{b}^\dagger \hat{b}), \quad [\hat{b}, \hat{b}^\dagger] = 1 \\ &= m \left(\frac{1}{2} + \hat{b}^\dagger \hat{b} \right), \quad : \hat{H} : \equiv m \hat{b}^\dagger \hat{b} \end{aligned} \quad (16)$$

↪ normal ordering (creation op's to left)
 $: \hat{b} \hat{b}^\dagger : \equiv \hat{b}^\dagger \hat{b}$

≡ omission of cc constants
 (choice of E scale)

$$\phi \rightarrow \hat{\phi} = \frac{1}{\sqrt{2m}} (\hat{b} + \hat{b}^\dagger)$$

↪ Heisenberg rep.

$$\begin{aligned} \hat{\phi}^H &= e^{i\hat{H}t} \hat{\phi} e^{-i\hat{H}t} = e^{im(\frac{1}{2} + \hat{b}^\dagger \hat{b})t} \frac{1}{\sqrt{2m}} (\hat{b} + \hat{b}^\dagger) e^{-im(\frac{1}{2} + \hat{b}^\dagger \hat{b})t} \\ &\quad \text{"-|n\rangle"} \rightarrow \left\{ e^{im(\frac{1}{2} + n)t} \frac{1}{\sqrt{2m}} \hat{b} + e^{-im(\frac{1}{2} + n)t} \frac{1}{\sqrt{2m}} \hat{b}^\dagger \right\} e^{-im(\frac{1}{2} + n)t} \\ &= \frac{1}{\sqrt{2m}} (\hat{b} e^{-int} + \hat{b}^\dagger e^{int}) \quad \text{"field operator" (17)} \end{aligned}$$

→ homework problems (1), (2)

⇒ 'recipe' for canonical quantization

- 1 $\mathcal{L} \xleftrightarrow{\partial_\mu \partial_{\partial_\mu \varphi} \mathcal{L} = \partial_\varphi \mathcal{L}} \text{field eqns.}$
 - 2 $\pi_i = \partial \dot{\varphi}_i \mathcal{L}, \quad \mathcal{H} = \pi_i \dot{\varphi}_i - \mathcal{L}$
 - 3 Bose $[\varphi_i(t, \vec{r}), \pi_j(t', \vec{r}')] = i\delta(t-t')\delta(\vec{r}-\vec{r}')\delta_{ij}$
Fermi $\{ \psi, \psi \} = \dots$
 - 4 expand φ, π in terms of solutions of field eqns.
coefficients re creation/annihilation op's
 $\varphi \rightarrow$ field operator
- (18)

some theories one should know

• harmonic oscillator: $\mathcal{L} = \frac{1}{2}\dot{\phi}^2 - \frac{1}{2}m^2\phi^2$ (1+1 D free field theory) (19)

• ϕ^4 theory: $\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4$, $\phi(x)$ (the toy-model, ϕ scalar, real re neutral particles) (20)

• QED: $\mathcal{L} = \bar{\psi}(i\not{D}-m)\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$ (electrons + photons) (21)

• QCD: $\mathcal{L} = \bar{\psi}(i\not{D}-m)\psi - \frac{1}{4}F_{\mu\nu}^a F^{\mu\nu a}$ (quarks + gluons) (22)

$A_\mu = A_\mu^a T^a, \quad F_{\mu\nu} = \frac{i}{2}[D_\mu, D_\nu] = \partial_\mu A_\nu - \partial_\nu A_\mu - ig[A_\mu, A_\nu] \equiv F_{\mu\nu}^a T^a$
 \uparrow generators of $SU(N)$, $a=1 \dots N^2-1$ (reality: $N=3$ colours \rightarrow 8 gluons)

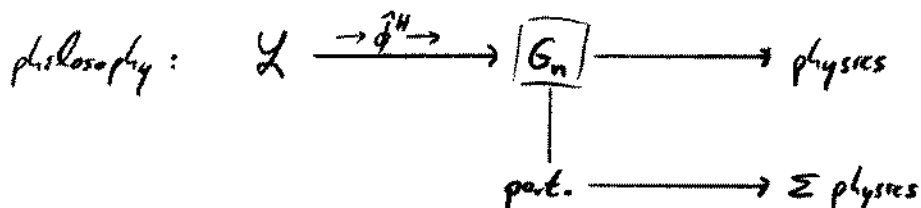
• Higgs model: $\mathcal{L} = (D_\mu \phi)^\dagger D^\mu \phi + \mu^2 \phi^\dagger \phi - \lambda(\phi^\dagger \phi)^2 - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$ (neutral/charged scalar + (SSB!) massive gauge bosons) (23)

• Standard Model (see page FT²/₉, chapter 2)

why be happy with $\hat{\phi}^H$?

→ def. Green's fcts ! time ordering: bigger times to the left

$$G_n(x_1, \dots, x_n) = \langle 0 | \mathcal{T} \hat{\phi}^H(x_1) \dots \hat{\phi}^H(x_n) | 0 \rangle \quad (24)$$



ex $G_2(t) = \langle 0 | \mathcal{T} \hat{\phi}(t) \hat{\phi}(0) | 0 \rangle$

((transl. inv.: $t_1, t_2 \rightarrow G_2(t_1 - t_2)$))

$$= \frac{1}{2m} \theta(t) \langle 0 | (b e^{-imt} + \cancel{b^\dagger e^{imt}}) (\cancel{b^\dagger} + b) | 0 \rangle$$

$$+ \frac{1}{2m} \theta(-t) \langle 0 | (\cancel{b} + b^\dagger) (\cancel{b} e^{-imt} + b^\dagger e^{imt}) | 0 \rangle$$

$$= \frac{1}{2m} (\theta(t) e^{-imt} + \theta(-t) e^{imt}) \langle 0 | b b^\dagger | 0 \rangle = 1$$

$$= \frac{1}{2m} e^{-im|t|} \quad (25a)$$

$$(\partial_t^2 + m^2) i G_2(t) = \delta(t) \quad (\text{that's the reason for the name})$$

Fourier transform ?

$$\tilde{G}_2(\omega) = \int dt e^{i\omega t} \frac{1}{2m} e^{-im|t|} \left(e^{-\epsilon|t|} \right) \quad \text{convergence factor. see below.}$$

$$= \frac{1}{2m} \left(\frac{e^{i\omega t}}{(i\omega + im + \epsilon)} \Big|_{t=-\infty}^0 + \frac{e^{i\omega t}}{(i\omega - im - \epsilon)} \Big|_0^{\infty} \right)$$

$$= \frac{1}{2m} \left(\frac{1}{i\omega + im + \epsilon} - \frac{1}{i\omega - im - \epsilon} \right) = \frac{i + \epsilon/m}{\omega^2 - (m - i\epsilon)^2}$$

$$= \frac{i}{\omega^2 - m^2 + i\epsilon} + O(\epsilon)$$

Feynman propagator
(in 0+1 D)

(25)

response function (example for "G_n → physics)

$$H_0 = H + H_{ext} = H - j(t)\phi \quad (L \rightarrow L + j\phi)$$

↑ some mean, unsolvable, many-body system ↑ source; perturbation

ask for $A(t) = \langle 0 | \hat{\phi}^H | 0 \rangle$

interaction picture wrt H_{ext}
 $|0, t=-\infty\rangle_0 = |0\rangle$ (H_{ext}(t=∞)=0)

$\langle \phi | \psi, t \rangle = -i H(t) | \psi, t \rangle$

$$= \langle 0 | (1 + \dots) \hat{\phi}^H(t) (1 - i \int_{-\infty}^t dt' H_{ext}(t')) | 0 \rangle$$

$$= -j(t') \hat{\phi}^H(t')$$

$$= i \int_{-\infty}^t dt' \langle 0 | [\hat{\phi}^H(t), \hat{\phi}^H(t')] | 0 \rangle j(t') + O(j^2)$$

$\underbrace{j(t')}_{= e^{\epsilon t'} j(t)}$ small!

$$e^{\epsilon t} a(t) =$$

from t'-integration limit

$$\Rightarrow a(t) = \int dt' \underbrace{\theta(t-t') e^{-\epsilon(t-t')}}_{\equiv \chi(t-t')} i \langle 0 | [\hat{\phi}^H(t), j(t')] | 0 \rangle$$

(26)

here (ham. oscillator) one can actually compute χ :

$$\chi(t) = \theta(t) e^{-\epsilon t} \frac{5m(mt)}{m}$$

causality

$$\tilde{\chi}(\omega) = -\frac{1}{(\omega + i\epsilon)^2 - m^2}$$

$$i \tilde{G}_2(\omega) = \tilde{\chi}(\omega) \text{ for } \text{Re } \omega > 0 \quad (G_n \rightarrow \text{physics!}) \quad (27)$$

outlook

Temperature? → next chapter.

generate all Greens, without canonical quant.

→ path integral! (T, T=0) at once; next chapter